



Integrability as a new method for exact results on quasinormal modes of black holes

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From gauge to gravity & back

- ▶ Recently, a surprising connection between (Nekrasov-Shatashvili (NS) deformed/quantum) $\mathcal{N} = 2$ supersymmetric (SUSY) $SU(2)$ gauge theories and black holes (BHs) perturbation theory emerged.
- ▶ It was first found that **quantization conditions on the gauge periods** a, A_D (building blocks of prepotential in quantum Seiberg-Witten theory)

$$A_D(a, \hbar, \Lambda_n) = i\pi n \quad n \in \mathbb{Z}, \quad (1)$$

allow to compute the quasinormal modes (QNMs) frequencies $\omega_n \sim \Lambda_n$ of BHs (of the gravitational waves signal in the ringdown phase). [arXiv:2006.06111]

- ▶ This result is important because:
 1. constitutes a new analytic characterization of QNMs;
 2. in the growing subsequent literature it allowed to derive many new results for the BHs theory;
 3. constitutes a **concrete, fruitful application of SUSY, beyond elementary particle physics.**

Integrability & N=2 gauge th.

- ▶ Using the ODE/IM correspondence (see below) **we connected basic integrability Q , T and Y functions to the quantum gauge theory periods a, a_D** . For instance we proved the basic identifications: [1]

$$\begin{aligned} Q(\theta, P) &= \exp \{2\pi i a_D(\hbar, u, \Lambda)/\hbar\} \\ T(\theta, P) &= 2 \cos \{2\pi a(\hbar, u, \Lambda)/\hbar\}. \end{aligned} \quad (2)$$

under the parameters correspondence

$$\frac{\hbar}{\Lambda} = \frac{\epsilon_1}{\Lambda} = e^{-\theta} \quad \frac{u}{\Lambda^2} = \frac{1}{2} P^2 e^{-2\theta}, \quad (3)$$

here for the self-dual Liouville Integrable Model (IM) and $SU(2)$ $N_f = 0$ gauge theory.

- ▶ This allowed also to **find several new results on both sides**, for instance:
 1. an exact non linear integral equation (Thermodynamic Bethe Ansatz, TBA) for the gauge periods;
 2. an interpretation of the integrability functional relations as new exact R -symmetry relations for the periods;
 3. new formulas for the local integrals of motion in terms of gauge periods.

For instance, the TQ relation

$$T(\theta, u) = \frac{Q(\theta - i\pi/2, -u) + Q(\theta + i\pi/2, -u)}{Q(\theta, u)} \quad (4)$$

turns out to be a new exact R -symmetry relation for the periods, reducing to the known asymptotic ones in the limit $\theta \rightarrow \infty$ (for the gauge periods expansion modes $a^{(n)}, a_D^{(n)}$)

$$a_D^{(n)}(-u) = i(-1)^n \left[-\text{sgn}(\Im u) a_D^{(n)}(u) + a^{(n)}(u) \right]. \quad (5)$$

- ▶ This is a **new kind of gauge/integrability correspondence** and we developed it until now for the $SU(2)$ $N_f = 0, 1, 2$ [1, 4] and $SU(3)$ $N_f = 0$ $\mathcal{N} = 2$ [2] (NS deformed) gauge theories, in correspondence with Liouville, Perturbed Hairpin, A_2 Toda IMs.

Integrability & quasinormal modes of black holes

- ▶ Perturbations Φ of a BHs are mathematically solutions of some PDE, which expands in **quasinormal modes (QNMs)** ω_n as

$$\Phi(t, x) = \sum_n e^{i\omega_n t} \text{Res} \left(\frac{1}{W(\omega)} \right) \Big|_{\omega_n} \int_{-\infty}^{\infty} \psi_{-}(\omega_n, x_{<}) \psi_{+}(\omega_n, x_{>}) \mathcal{I}(\omega_n, x') dx', \quad (6)$$

where x is the tortoise coordinate, $\mathcal{I}$ is given by initial data on the perturbation and $\psi_{\pm}$ are solutions of the ODE associated to the PDE through Laplace transform.

- ▶ Crucially, we noticed that **QNMs are given by the zeros of the same wronskian** which in ODE/IM correspondence (see below) defines **the Baxter's Q function** of some IM [3]

$$W(\omega_n) = W[\psi_{+}, \psi_{-}](\omega_n) = Q(\omega_n) = 0. \quad (7)$$

In particular, ODE (14) for the self-dual Liouville IM describes the scalar perturbation of the D3 brane

$$ds^2 = H(r)^{-1/2} (-dt^2 + d\mathbf{x}^2) + H(r)^{1/2} (dr^2 + r^2 d\Omega_5^2), \quad H(r) = 1 + L^4/r^4, \quad (8)$$

under the change of parameters

$$r = L e^{\frac{y}{2}} \quad \omega L = -2ie^{\theta} \quad P = \frac{1}{2}(l+2), \quad (9)$$

- ▶ Thanks to our identifications (2) **we prove that QNMs are given by quantization conditions on gauge period a_D**

$$2\pi i a_D(i\hbar, -u, \Lambda_n) = \pi \left(n + \frac{1}{2} \right), \quad n \in \mathbb{N} \quad (10)$$

and prove the result (1), using also consistent relations between gauge periods a_D, A_D .

$$Q(\theta, P) = \exp \{2\pi i a_D(\hbar, u, \Lambda)/\hbar\} = i \frac{\sinh A_D(\hbar, a, \Lambda)/\hbar}{\sinh 2\pi i a/\hbar}, \quad (11)$$

- ▶ It follows also a new simple and effective method to **numerically compute the quasinormal modes: the Thermodynamic Bethe Ansatz (TBA)** for $\varepsilon = -2 \ln Q$

$$\varepsilon(\theta) = \frac{16\sqrt{\pi^3}}{\Gamma(\frac{1}{4})^2} e^{\theta} - 2 \int_{-\infty}^{\infty} \frac{\ln[1 + \exp\{-\varepsilon(\theta')\}]}{\cosh(\theta - \theta')} \frac{d\theta'}{2\pi}, \quad (12)$$

which we compare with other standard methods and sometimes find convenient.

n	l	TBA	Leaver	WKB approx.
0	0	<u>1.36912</u> - <u>0.504048i</u>	<u>1.36972</u> - <u>0.504311i</u>	1.41421 - 0.5i
0	1	<u>2.09118</u> - <u>0.501788i</u>	<u>2.09176</u> - <u>0.501811i</u>	2.12132 - 0.5i
0	2	<u>2.8057</u> - <u>0.501009i</u>	<u>2.80629</u> - <u>0.501000i</u>	2.82843 - 0.5i
0	3	<u>3.51723</u> - <u>0.500649i</u>	<u>3.51783</u> - <u>0.500634i</u>	3.53553 - 0.5i
0	4	<u>4.22728</u> - <u>0.500453i</u>	<u>4.22790</u> - <u>0.500438i</u>	4.24264 - 0.5i

- ▶ The spacetimes for which we show this are in the simplest $SU(2)$ $N_f = 0$ case the **D3 brane** in the $N_f = 1, 2$ case a **generalization of extremal Reissner-Nordström (charged) BHs** [3, 4].
- ▶ Anyway, the ODE/IM construction can be realized still for the $SU(2)$ $N_f = 3, 4$ theories, which correspond to non-extremal, asymptotically flat and asymptotically AdS BHs.

ODE/IM correspondence

- ▶ In this classic approach to integrability, the Q function of some Integrable Model (IM) is defined from the regular solutions at different singular points

$$Q = W[\psi_{+}, \psi_{-}] \quad \psi_{\pm}(y) \rightarrow 0 \quad y \rightarrow \pm\infty \quad (13)$$

of some ODE, like (self-dual Liouville IM) [1]

$$-\frac{d^2}{dy^2} \psi + [e^{2\theta}(e^y + e^{-y}) + P^2] \psi = 0. \quad (14)$$

- ▶ One can derive also T, Y functions as well as their functional and integral equations they satisfy.
- ▶ It **allows to apply integrability to other different physical theories** which share the ODE.

Possible future developments

- ▶ Besides **models extension**, it is very intriguing to study the **applications of integrability beyond QNMs** (e.g. the **greybody factor** seems a ratio of Qs.).
- ▶ Studying AdS BHs ($SU(2)$ $N_f = 4$ gauge theory) would allow to extend the triple Integrability/Gauge/Gravity correspondence to a **4-fold correspondence** through **holography**.

References

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- [2] D. Fioravanti, H. Poghosyan, R. Poghosyan, arXiv:1909.11100
- [3] D. Fioravanti, D. Gregori, arXiv:2112.11434
- [4] D. Fioravanti, D. Gregori, H. Shu, to appear